Detector Basics II

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Neutrino Summer School Fermilab, July 7th, 2009

EXO	/ 1
CUORE	// 2
Majorana/GERDA	/// 3
SuperNEMO	//// 4
KATRIN	// 2
LVD	/ 1
KamLAND	/ 1
BOREXINO	// 2
DoubleCHOOZ	////// 7
Daya Bay	/////// 8
Hanohano	/1
SNO+	// 2
MiniBooNE	// 2
MicroBooNE	// 2
MINOS	///////// 10
MINERVA	// 2
T2K/SK/MEMPHYSIS	///////// 10
NOVA	///// 6
LqAr	//// 4
Ice Cube	//// 4
Pierre Auger	/ 1
	1

Results of my index card survey

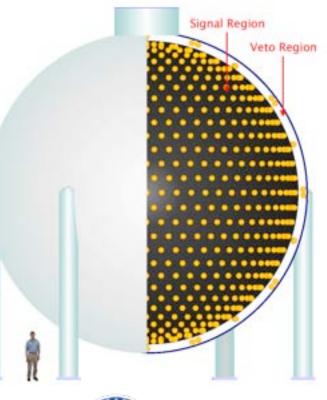
Answers to questions posted at: http://enrico1.physics.indiana.edu/messier/post/nss09qanda.txt

Neutrino detectors

- Topics for the remaining two lectures
- Today
 - Cherenkov detectors
 - Tracking calorimeters
- Thursday
 - tau neutrino detection
 - Large liquid scintillator detectors
 - Time projection chambers

Cherenkov detectors





MiniBooNE Detector

SNO

6000 mwe overburden

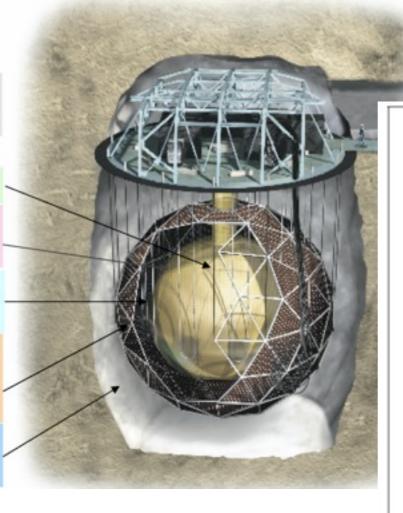
1000 tonnes D₂O

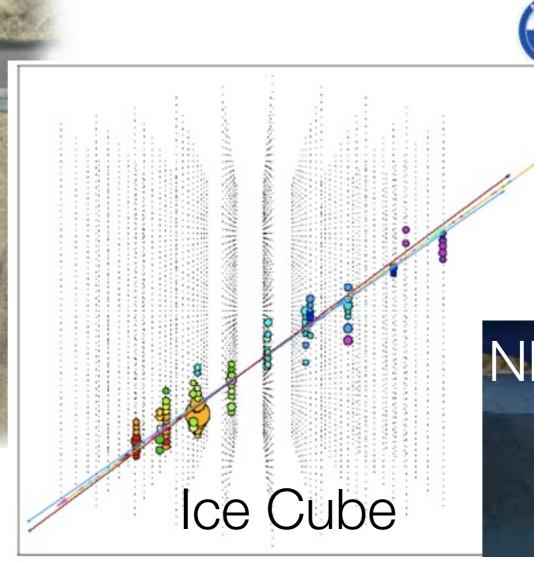
12 m Diameter Acrylic Vessel

1700 tonnes Inner Shield H₂O

Support Structure for 9500 PMTs, 60% coverage

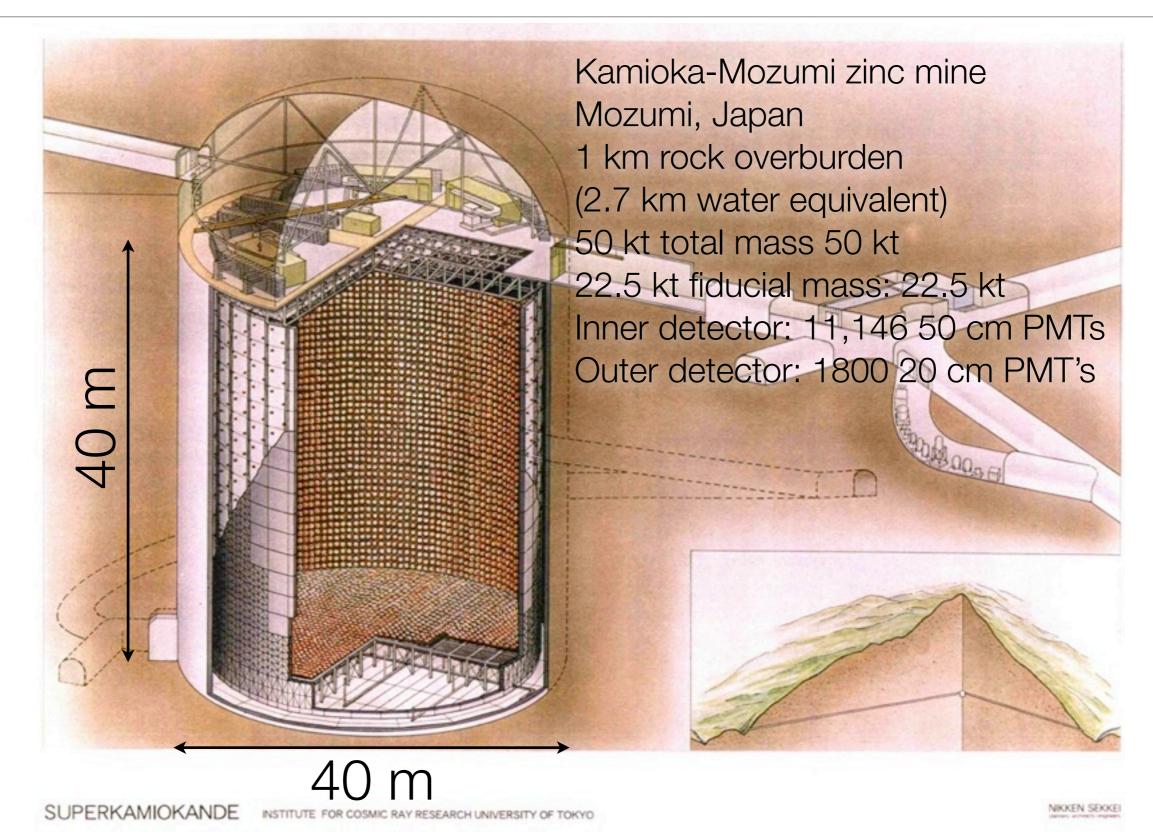
5300 tonnes Outer Shield H₂O





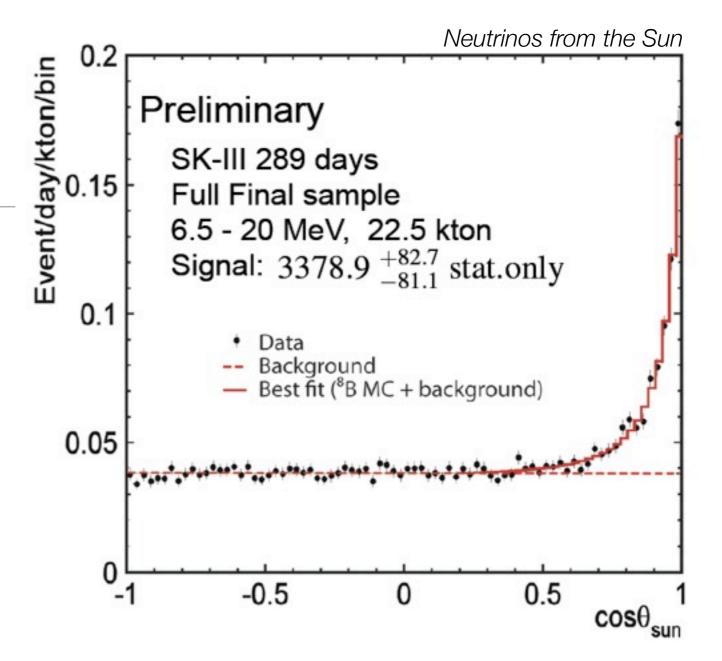


Super-Kamiokande



General performance

- Sensitive to a wide range of energies. Capable of electron and photo detection down to ~5 MeV
- Tracks produce rings on the walls.
 In high multiplicity events overlap of rings makes reconstruction difficult. Typically, analyses focus on quasi-elastic events which are very often single-track events.
- For single track QE events, neutrino energy reconstructed from kinematics (see next slides)



 Events with pions (and other tracks) that are below Cherenkov threshold lead to backgrounds for the quasi-elastic selection

SNO

6000 mwe overburden

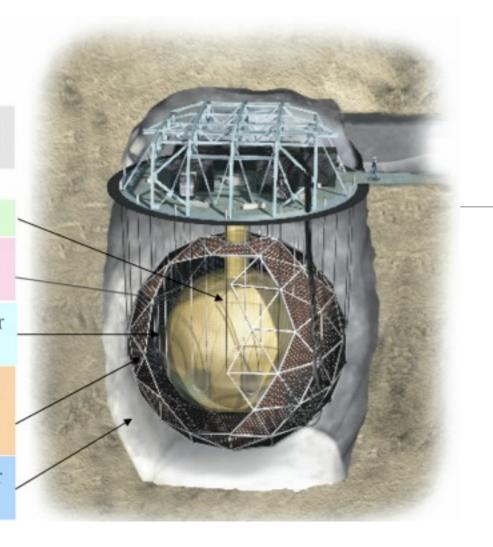
1000 tonnes D2O

12 m Diameter Acrylic Vessel

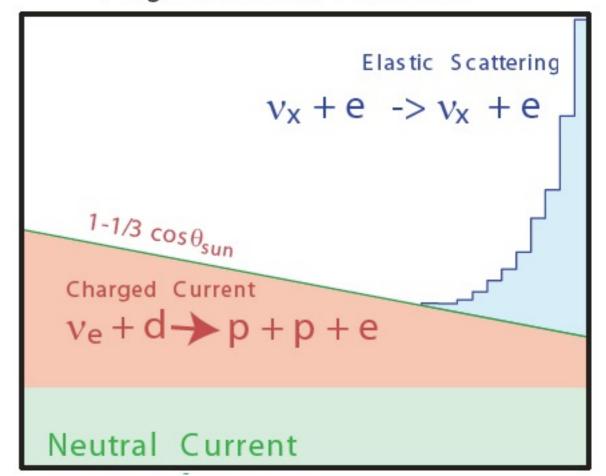
1700 tonnes Inner Shield H₂O

Support Structure for 9500 PMTs, 60% coverage

5300 tonnes Outer Shield H₂O



Angle between v and sun



• SNO detector used 1 kt D₂O instead of ordinary water. Provided additional detection channels at low energy: $\nu_x + e^- \rightarrow \nu_x + e^- \rightarrow \nu_x + e^- \rightarrow \nu_x + e^-$

$$\nu_e + d \rightarrow p + p + e^-$$
 CC

$$\nu_x + d \rightarrow p + n + \nu_x$$
 NC

$$n+d \rightarrow t + \gamma + 6.25 \text{ MeV}$$

• Neutron tagging by: $n+{}^{35}Cl \to {}^{36}Cl+\gamma+8.6~{
m MeV}$ $n+{}^3He \to p+t+0.76~{
m MeV}$

Cherenkov effect

 If speed of charged particle exceeds speed of light in a dielectric medium of index of refraction n, a "shock wave" of radiation develops at a critical angle:

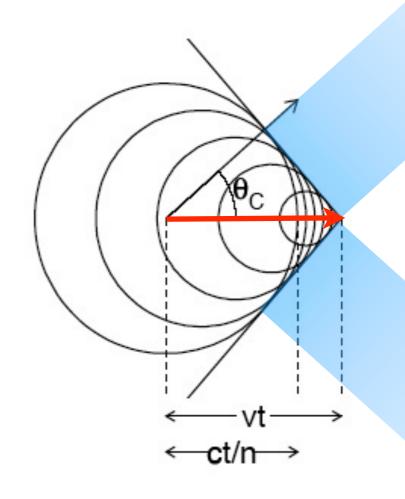
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 $\cos \theta_C = \frac{1}{\beta n}, \beta > \frac{1}{n}$

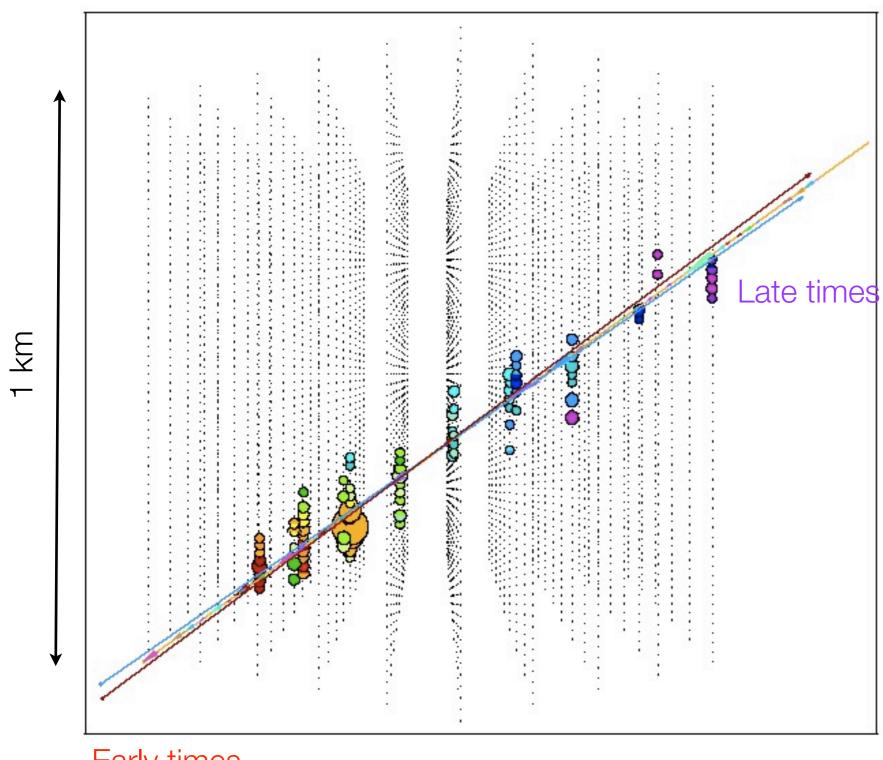
 PMTs record time and charge which provide unique solution for track position and direction. For Nhit PMTs measuring light arrival time t, minimize:

$$\chi^2 = \sum_{i=1}^{N_{\text{hit}}} \frac{(t_i - TOF_i)^2}{\sigma_t^2}$$

where TOF is the time of flight for photons to go from the track to the PMT



10 TeV neutrino induced muon neutrino in Ice Cube



Times differ by roughly 2.5 usec. For PMT with ~10 ns time resolution this gives an up vs. down discrimination of > 250sigma!

Early times

Cherenkov effect

Threshold means that slow particles produce no light. As particles come to a stop their rings collapse. Useful for particle ID near threshold.

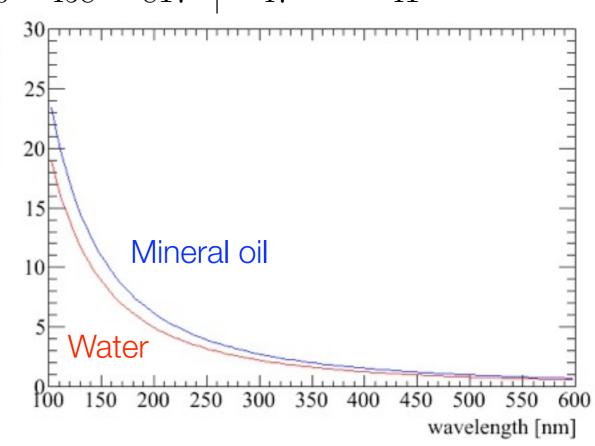
$$p_{\text{thresh}} = m\sqrt{\frac{1}{n^2 - 1}}$$

		$p_{\rm thresh} \ [{ m MeV/c}]$			$ heta_C$			
		e	μ	π	K	p	$\beta = 1$	$\beta = 0.9$
Water	n = 1.33	0.58	120	159	563	1070	42	33
Mineral Oil	n = 1.46	0.47	98	130	458	817	47	41

Number of photons produced per unit path

Number of photons produced per unit path length:
$$\frac{d^2N}{dxd\lambda} = \frac{2\pi\alpha z^2}{\lambda^2} \left(1 - \frac{1}{\beta^2 n(\lambda)^2}\right) = 370\sin^2\theta_C(E)/\text{eV/cm}$$
In both oil and water the useful part of this spectrum is between 300 and 600 pm

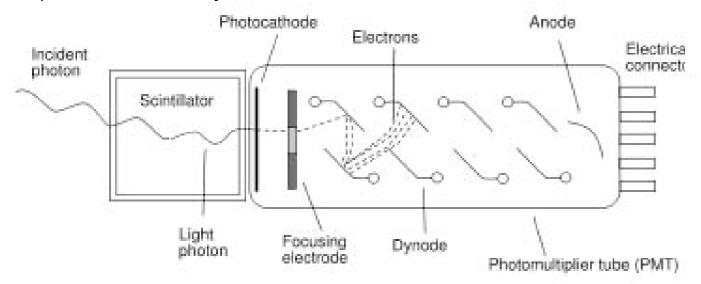
spectrum is between 300 and 600 nm bracketed by Rayleigh scattering on the low end and absorption on the high end



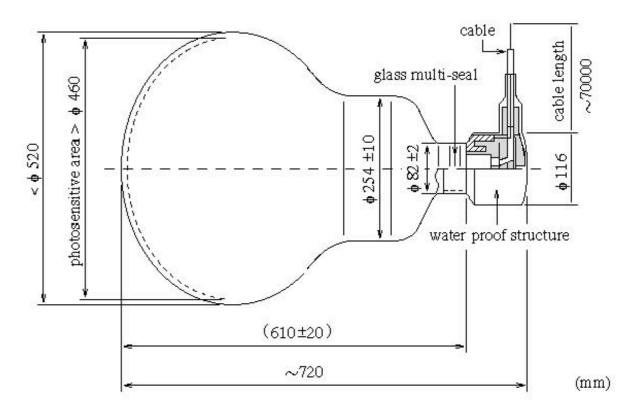
Photomultiplier tubes

Photon incident on the photocathode produces a photoelectron via the photoelectric effect. Probability to produce a photoelectron is called the quantum efficiency of the PMT.

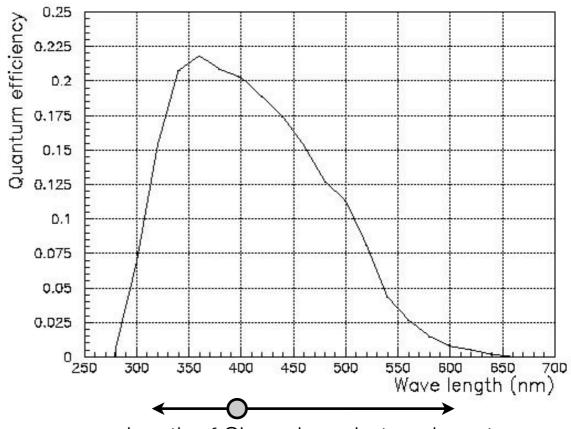
Output signal is seen as a current delivered to the *anode*. Typical gains are 10⁶ yielding pC-scale currents



A series of plates called *dynodes* are held at high voltage by the *base* such that electrons are accelerated from one dynode to the next. At each stage the number of electrons increases. Probability to get first electron from the photocathode to the first dynode is called the *collection efficiency*.



100 ns transit time, 2.2 ns time resolution



wavelength of Cherenkov photons in water

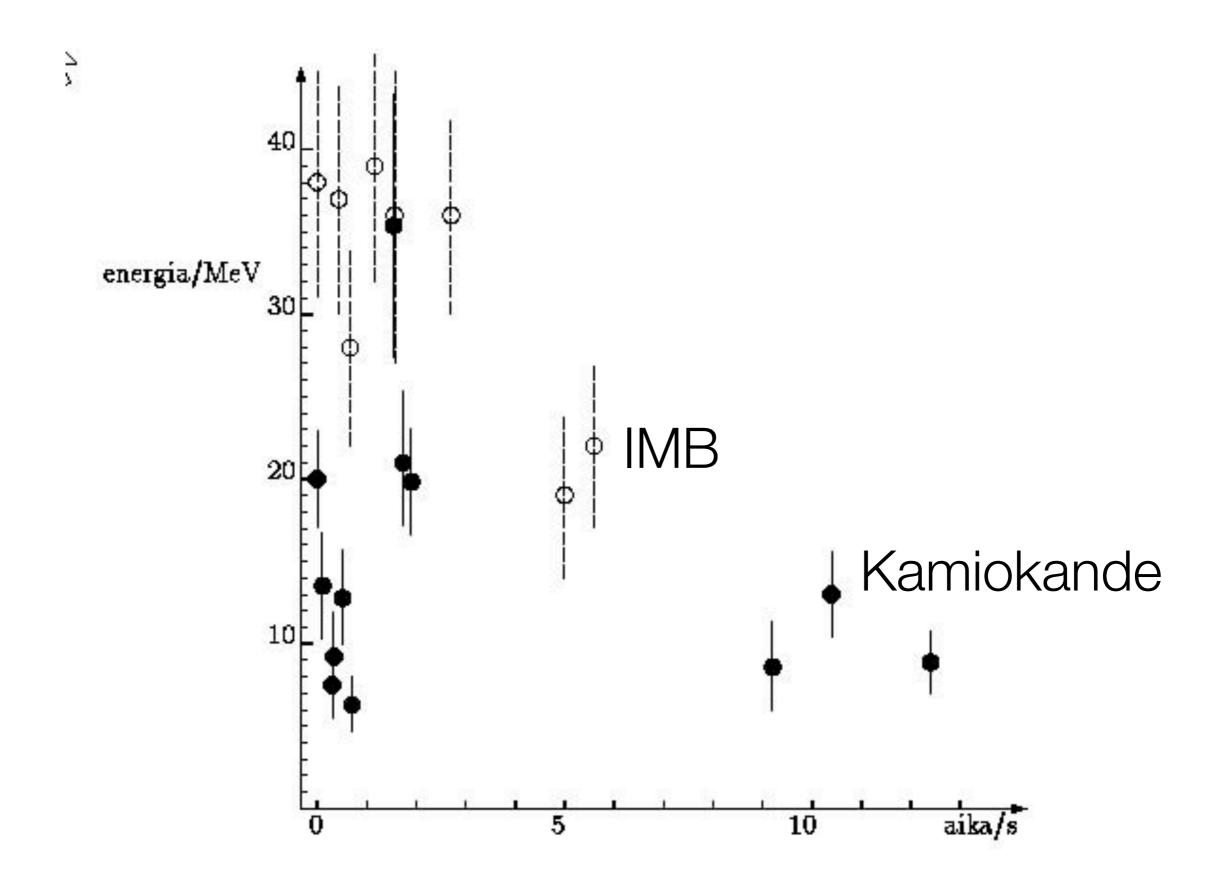
Quiz

Q: Estimate the vertex resolution for a water Cherenkov detector for a 10 MeV electron produced by the elastic scatter of a solar neutrino. Assume 40% of the detector walls are covered by PMT's and that the PMT's have an average of 25% efficiency. Estimate the energy resolution at this energy.

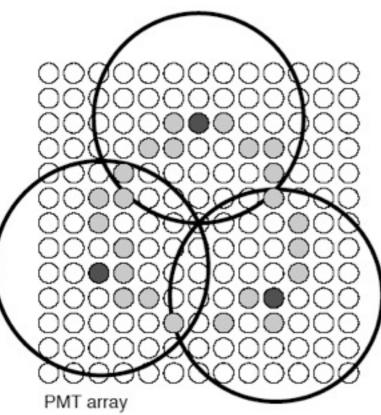
A: A 10 MeV electron will go about 5 cm in the tank making about N = $370*\sin(42^\circ)^2 = 160$ photons. Of those (0.4*0.25) = 0.1 will be detected. So I have ~16 detected photons each with a timing resolution of 2 ns ~= (60 cm * 1.33) = 80 cm since the speed of light is n * 30 cm/ns. This gives a final resolution of about: 80 cm/sqrt(16) = 20 cm. Energy resolution dominated by Poisson fluctuations on the number of photons collected. In this case ~sqrt(16)/16 = 25%.

Q: Compare the detection efficiencies for the Kamiokande (20% photocathode coverage) and IMB-1 (1% photocathode coverage) for a 15 MeV super-nova neutrino

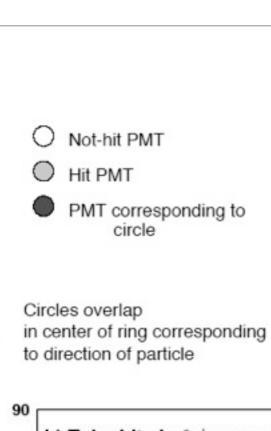
A: 15 MeV corresponds to about 240 photons which is about 0.6 detected photons on average in IMB and 12 in Kamiokande. Efficiency for detection is roughly $1-\exp(-0.6) = 45\%$ for IMB and $1-\exp(-12) = 99.99\%$ for Kamiokande.



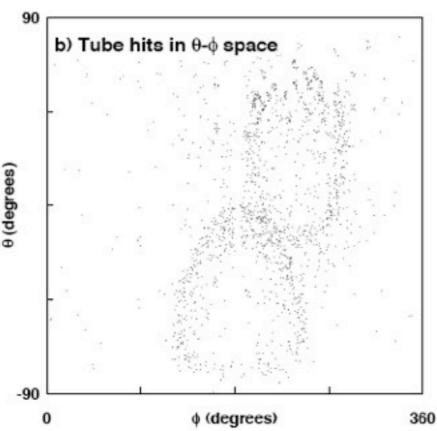
Water Cherenkov: Ring Counting

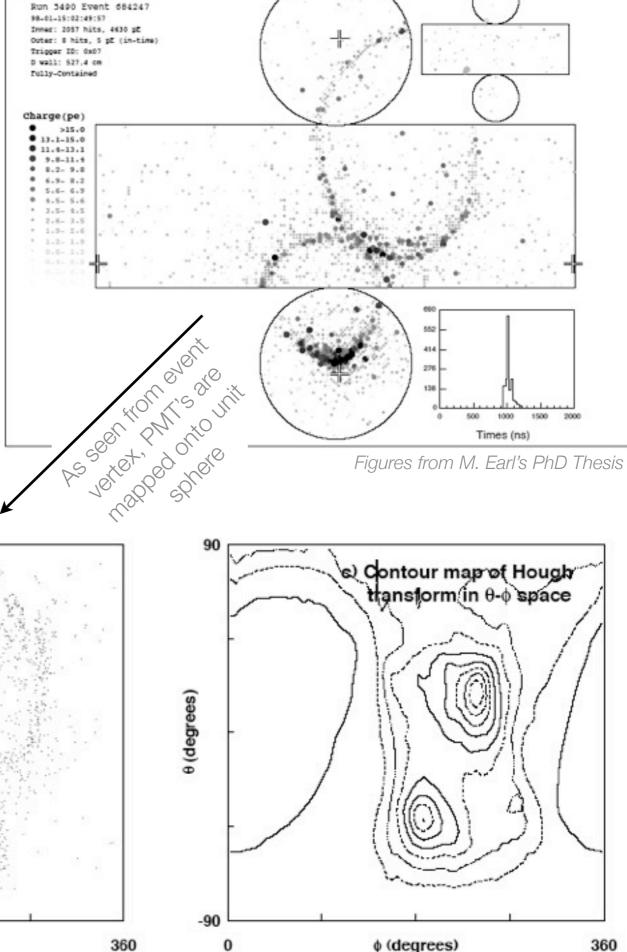


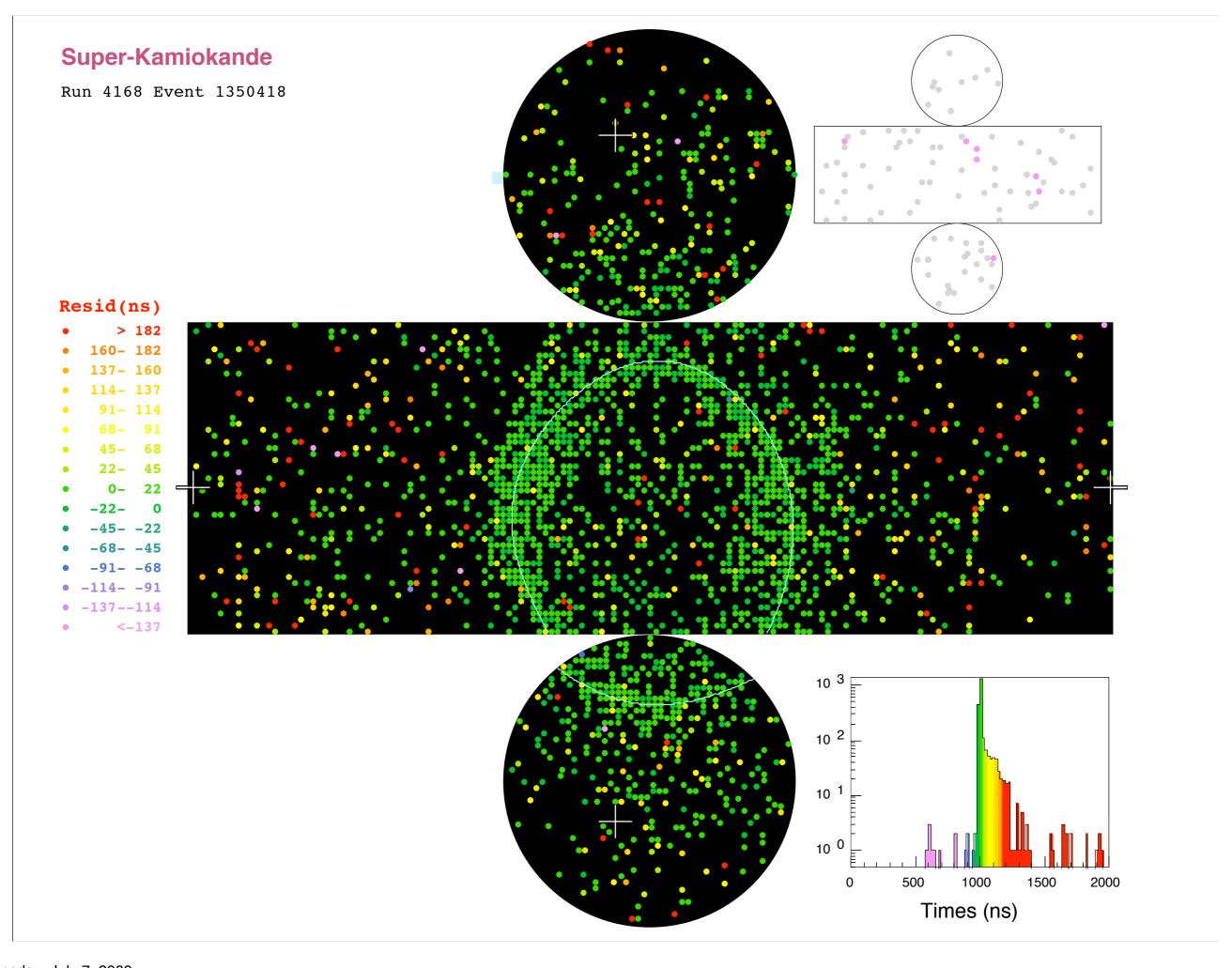
If you know the pattern you are looking for (line, circle, oval, etc.) the Hough transform is a method for converting a pattern recognition problem to a peak finding problem



Super-Kamiokande







Quasi-elastic reconstruction $\nu_{\mu} + n \rightarrow \mu^{-} + p$

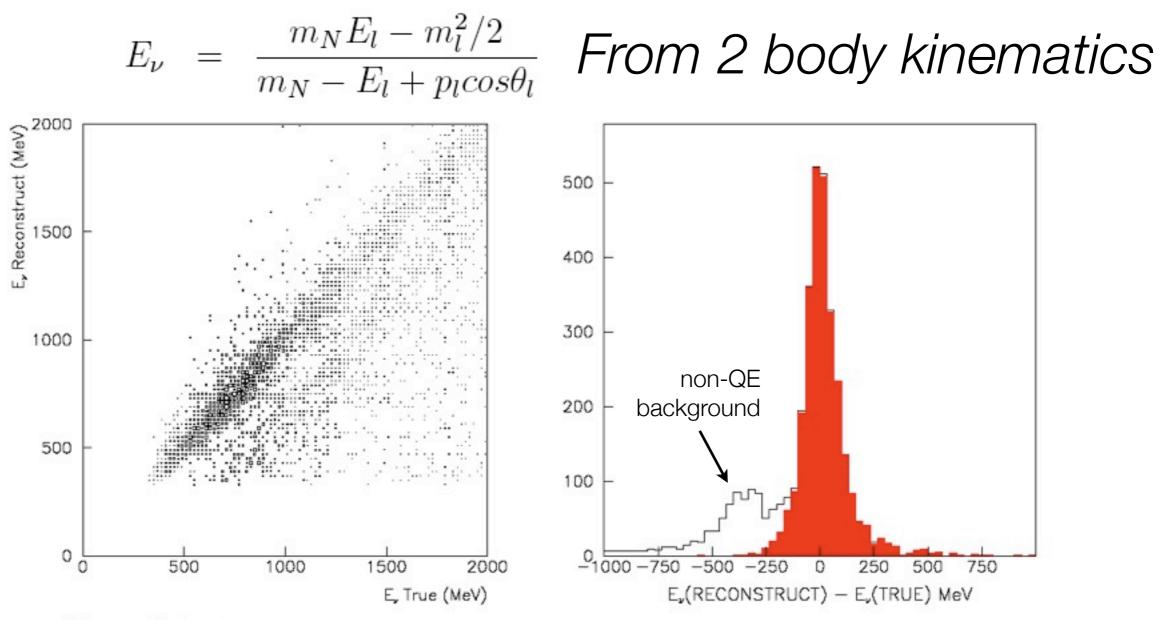
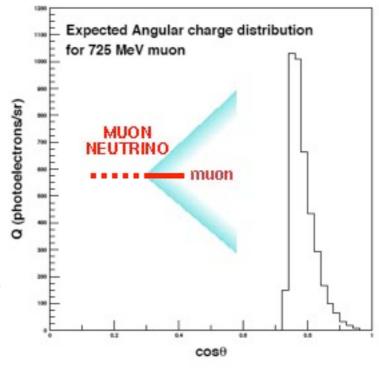


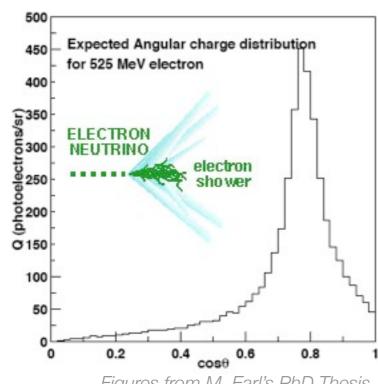
Figure 2: (left) The scatter plots of the reconstructed neutrino energy versus the true one for ν_{μ} events. The method of the energy reconstruction is expressed in Equation 14. (right) The energy resolution of ν_{μ} events for 2 degree off-axis beam. The shaded (red) histogram is for the true QE events.

Water Cherenkov: e/µ identification

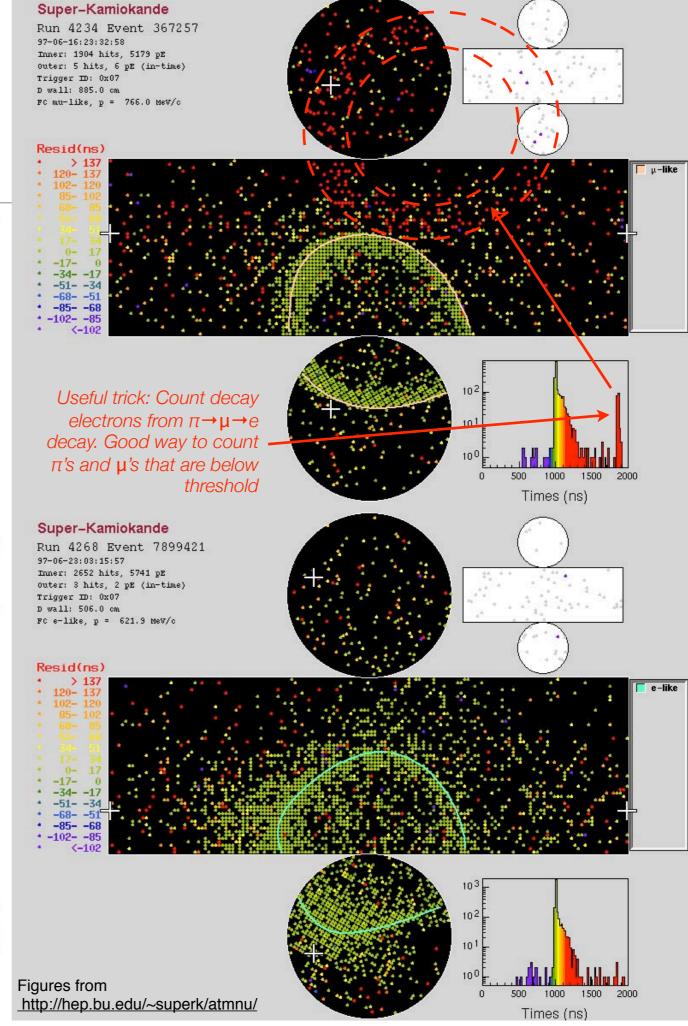
 At low momenta one can correlate the particle visible energy with the Cherenkov angle. Muons will have "collapsed" rings while electrons are ~always at 42°.



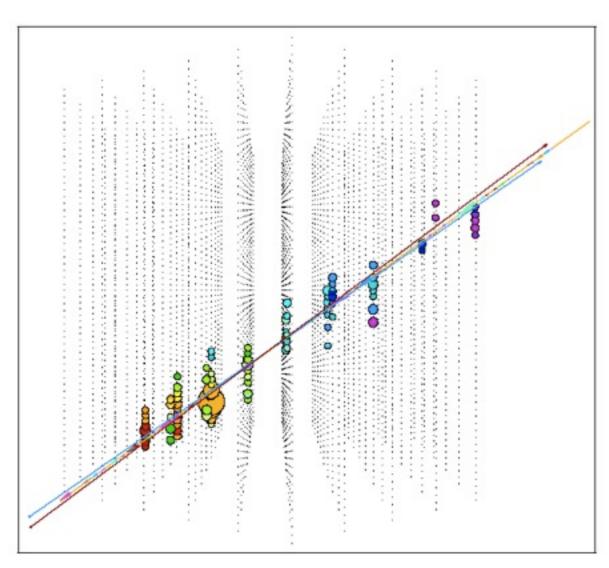
At higher momenta, look at the distribution of light around Cherenkov angle. Muons are "crisp", electron showers are "fuzzy". See plots and figures at the right.



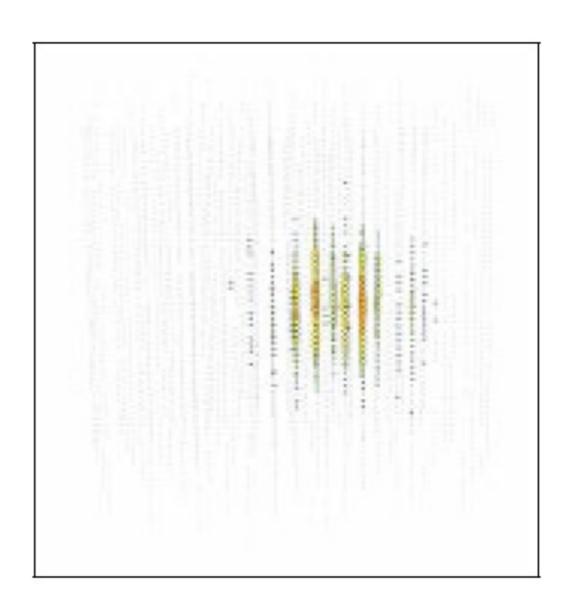
Figures from M. Earl's PhD Thesis



Particle ID in Ice Cube

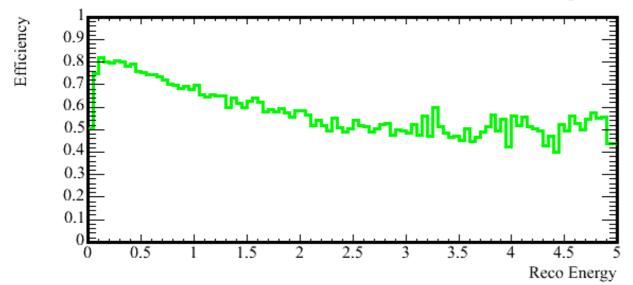


10 TeV muon neutrino induced upward muon



375 TeV electron neutrino

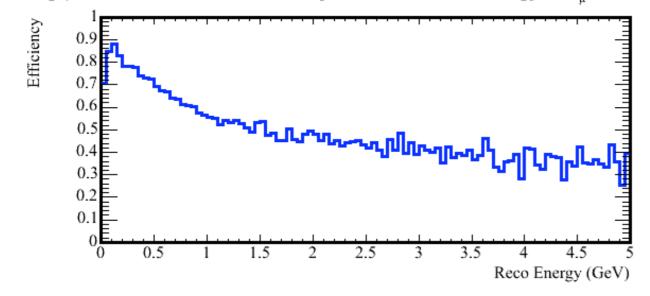
1-Ring, e-Like Reconstruction Efficiency vs Reconstructed Energy for $v_{\rm e}$ CC Events



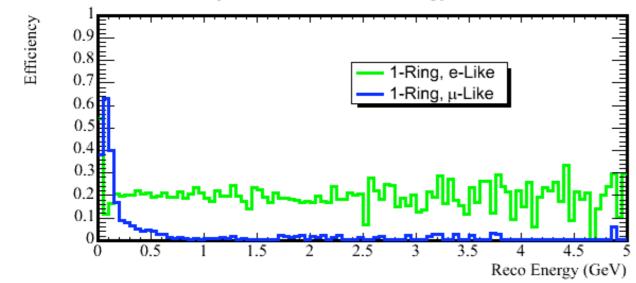
Additional selections:

no decay electrons: $\frac{\text{CC }\nu_{\mu}\,\text{NC }\text{CC }\nu_{e}}{14\%\,\,19\%\,\,76\%}$ signal energy window (T2K) 1% 16% 58% π^{0} likelihood fit 0.4% 10% 42%

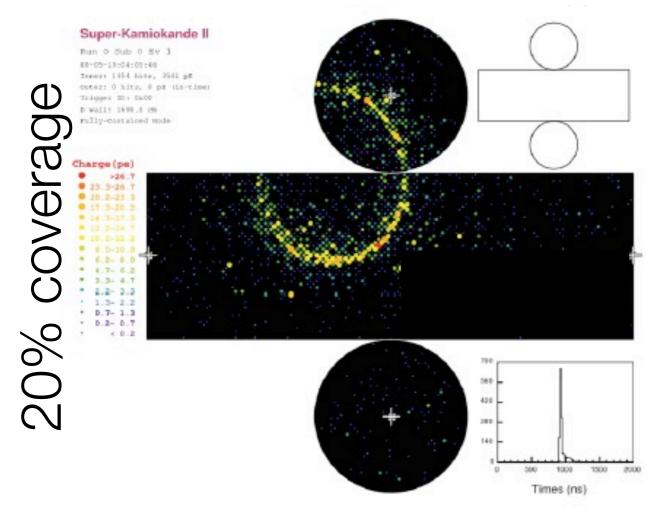
1-Ring, μ -Like Reconstruction Efficiency vs Reconstructed Energy for v_μ CC Events

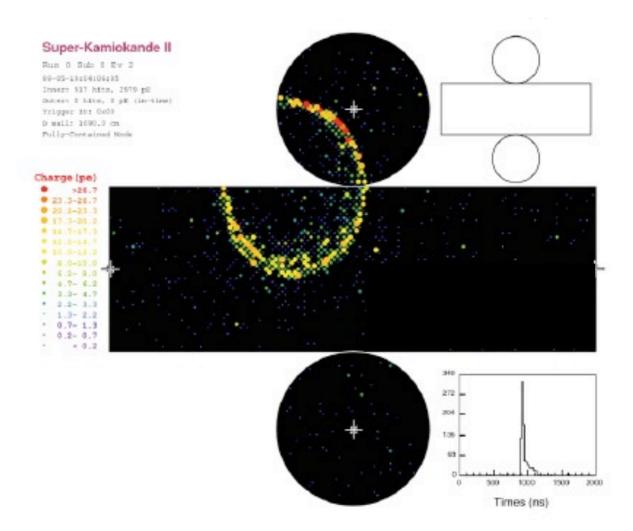


Reconstruction Efficiency vs Reconstructed Energy for NC Events

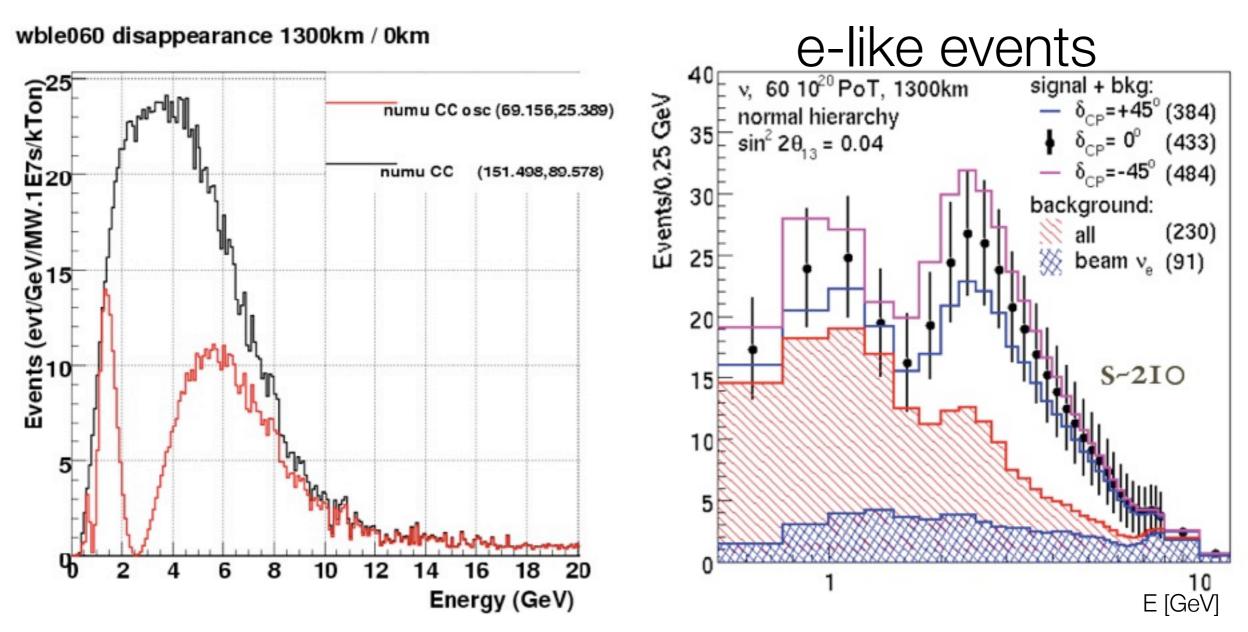


Notice: NC events much more likely to be e-like than μ -like due to π^0 production





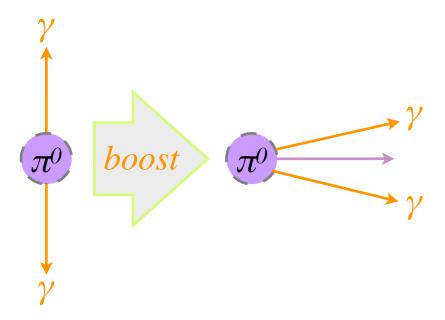
Pushing the technology: Sub-GeV to Multi-GeV

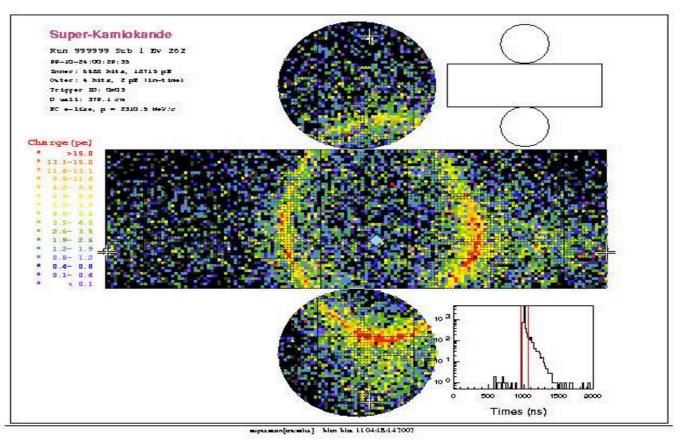


100 kt water detector in multi-GeV 2 MW wide band beam Fermilab to Homestake

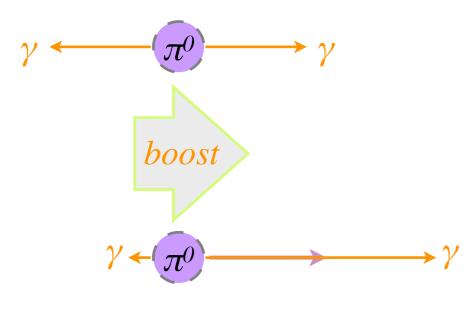
2 GeV visible energy One is signal, the other background

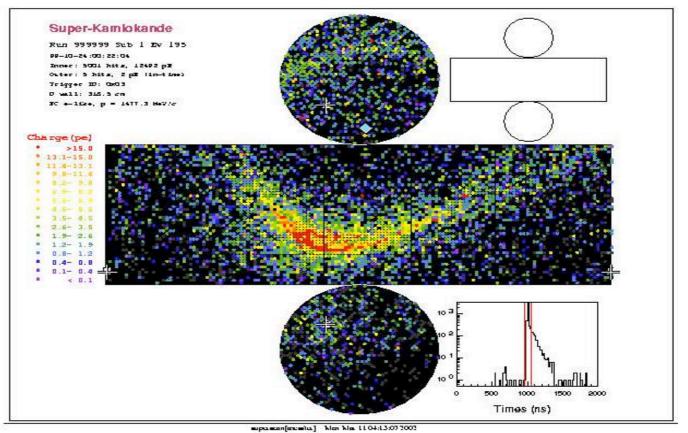
 π^0 decay at high energy





ve CC





 $NC \pi^0$

US DUSSEL Underground lab

4850 Level

Hyper-Kamiokande

Plat form
Liner
Opaque Sheet
Outer Detector
Inner Detector
Inner Detector
Inner Detector
Inner Detector
Inner Access Drift

SECTION

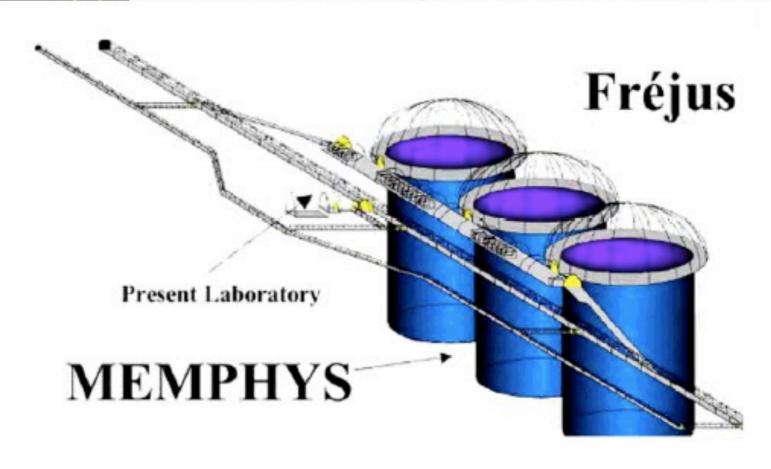
Photomultipliers

Comparison James Langer James L

(STRA

Mton fiducial volume: Total Length 800m (16 Compartments)

Megaton Scale
Water
Cherenkov
Detectors



20% or 40% Photocathode coverage?

PMT's cost ~\$3K USD and are one of the schedule drivers for construction of very large water Cherenkov detectors. Can you live with fewer?

	Super-K I (40% coverage)	Super-K II (20% coverage)
Sub-GeV vertex resolution	26 cm (e-like) / 23 cm (μ-like)	30 cm (e-like) / 29 cm (μ-like)
Sub-GeV particle mis-ID	0.81% (e-like) / 0.70% (μ-like)	0.69% (e-like) / 0.96% (μ-like)
Sub-GeV momentum resolution	4.8% (e-like) / 2.5% (μ-like)	6.3% (e-like) / 4.0% (μ-like)
p→e⁺πº signal efficiency	40.8±1.2 ±6.1%	42.2±1.2 ±6.3%
p→ e⁺πº background	0.39(±35%) events/100kty	0 events/100kty
p→ K⁺ν, γ tag signal efficiency	8.4±0.1 ±1.7%	4.7±0.1 ±1.0%
p→ K⁺ν, γ tag background	0.72(±28%) events/100kty	1.4(±30%) events/100kty
p→ K⁺ν, π⁺π⁰ signal efficiency	5.5±0.1 ±0.7%	5.7±0.1 ±0.4%
p→ K⁺ν, π⁺π⁰ background	0.59(±28%) events/100kty	1.0(±30%) events/100kty
T2K CCv _e likelihood effic.	83.7% (±0.1% stat)	84.8 %
T2K BG likelihood effic.	21.3 %	21.5 %

S.T. Clark, Ph.D. Thesis (2006) F. Dufour, T2KK Workshop (2006) <u>Preliminary</u> numbers, for comparison purposes. Final published efficiencies and BG may differ.

INO Tracking calorimeters 12m 150 100 50 -50 -100 SciBooNE 67 m -150 100 150 200 15.7 m NOvA 15.7 m

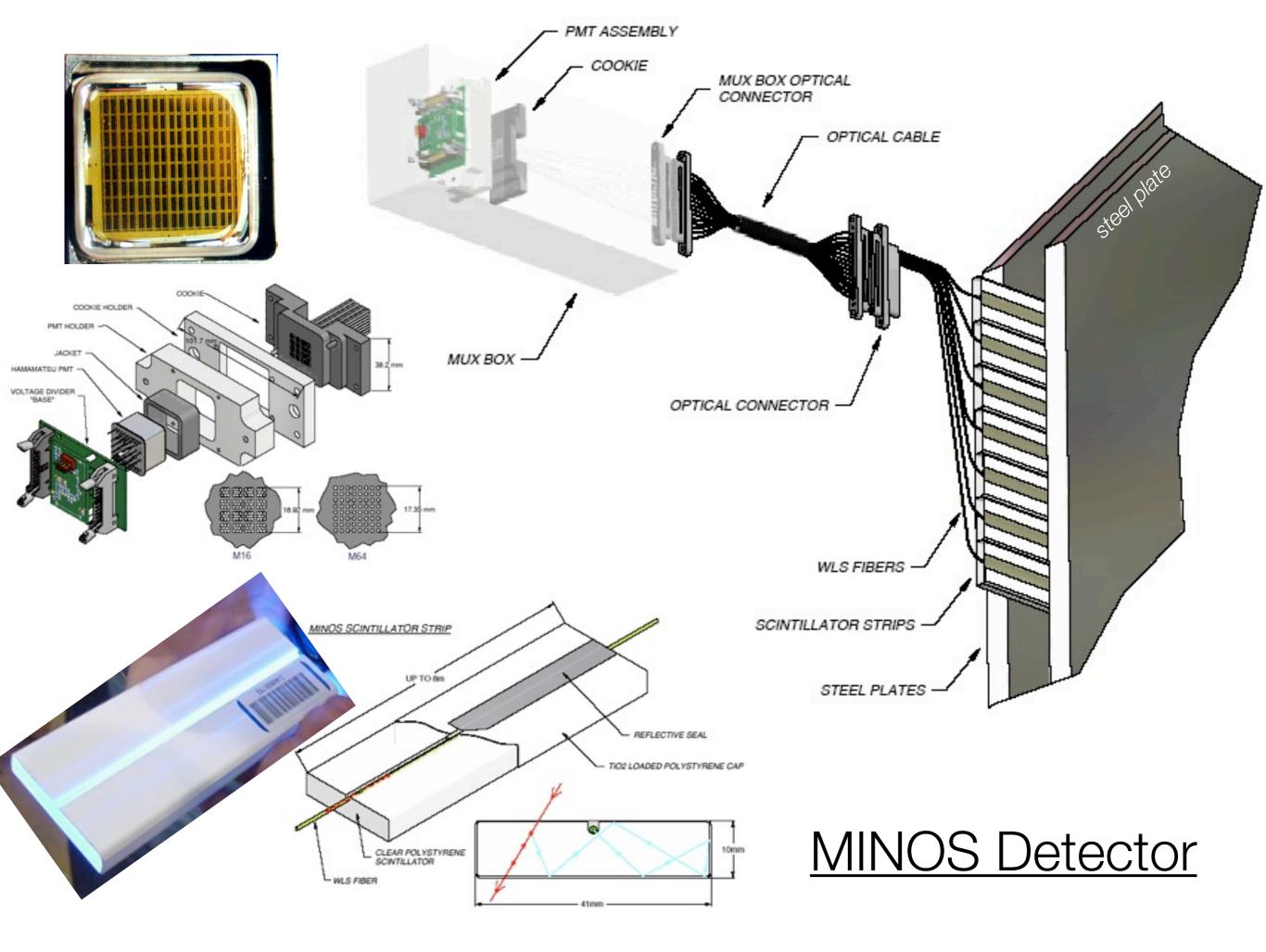
The MINOS Detectors

MINOS uses two functionally equivalent detectors:

- 2.54 thick magnetized steel plates
- 4.1 x 1 cm co-extruded scintillator strips
- optical fiber readout to multi-anode PMT's







scintillator modules layered on steel plane



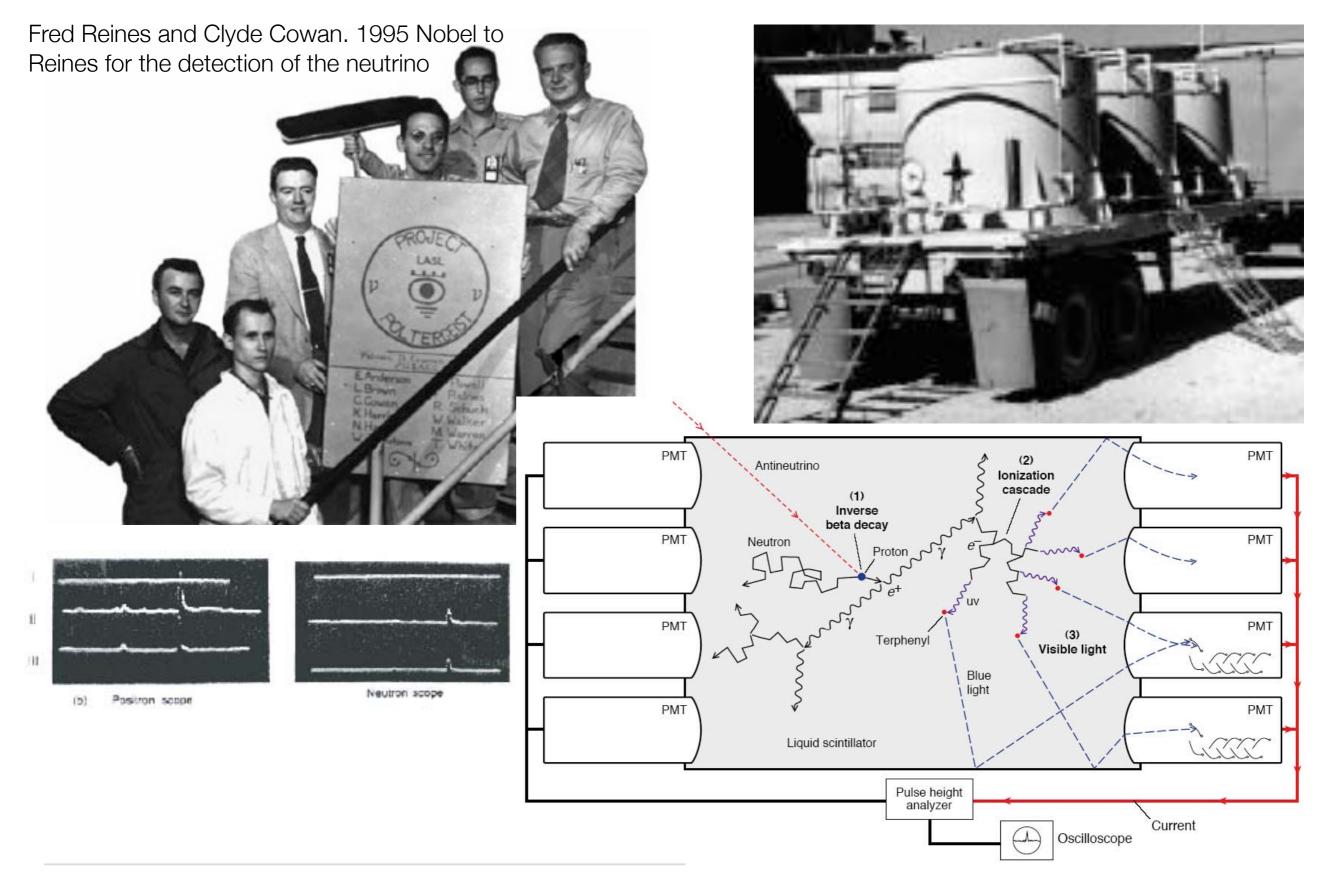
"strong back".
Removed after plane is hung in place

Scintillation process

- Scintillators are solid or liquid materials that produce light shortly after absorbing energy from a passing particle
- "Shortly" here is characterized by the decay time of the scintillator with the number of photons emitted varying as $n(t) = k(1 e^{-t/\tau})$. The fastest scintillators have tau's at low as 5 ns. Typical values are 10-100 ns.
- The number of scintillation photons produced per unit energy deposited goes as: $n = n_0 \frac{dE/dx}{1 + BdE/dx}$

where B is "Birk's" constant and accounts for saturation of the scintillator at high energy depositions.

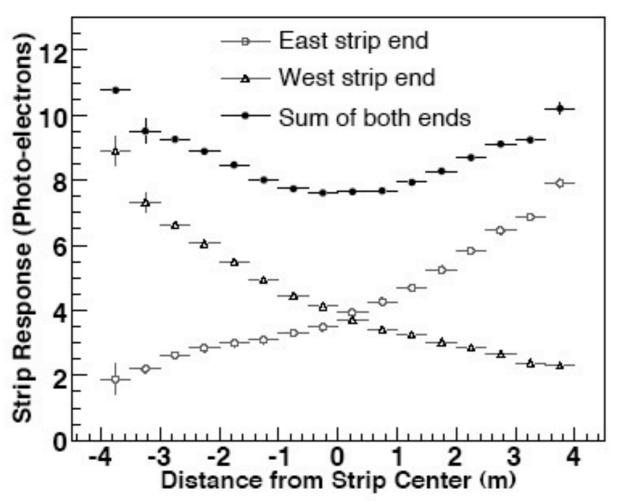
- Scintillation light is emitted in a distribution peaked typically around 350-400 nm. It is common to use compounds (eg. PPO, POPOP) which absorb this light and re-emit it at longer wavelengths where the scintillator has less absorption and where the fiber absorbs strongly.
- Light is captured by the fiber at typically 420 nm and reemitted at around 470 nm and is carried to the ends by total internal reflection. Transport characterized by a short attenuation length (~2 m) and a long attenuation length (~8 m).
- Final photon spectrum is well matched to wavelength response of PMT's



Los Alamos Science, Number 25 1997

Project Poltergeist, 1953

MINOS scintillator system



Single strip muon hit efficiency

Single sided:

$$\varepsilon = 1 - \exp(-2) = 86\%$$

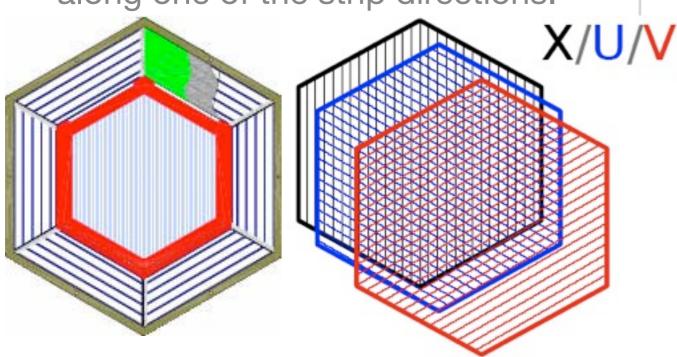
Double sided:

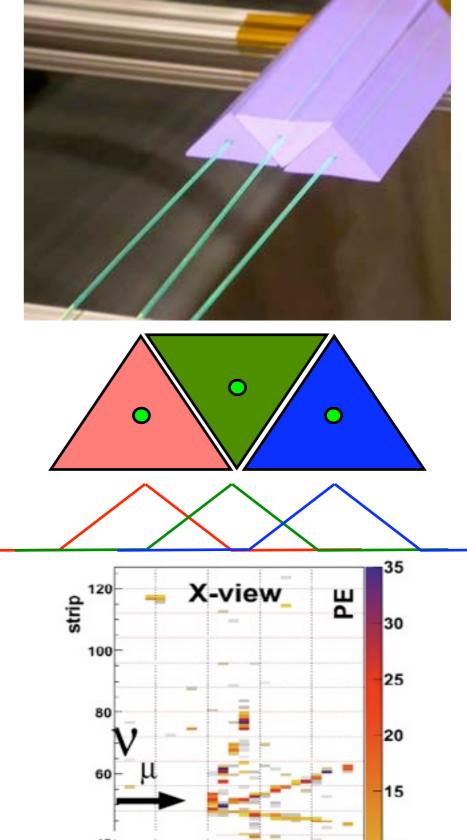
$$\varepsilon = 1 - \exp(-8) = 99.97\%$$

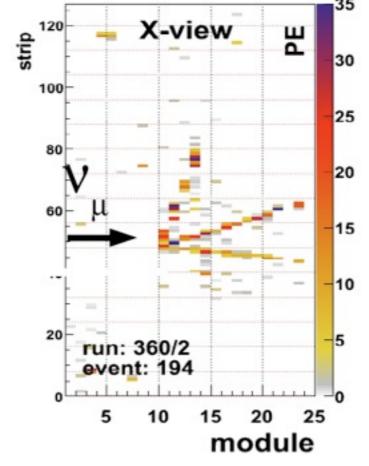
Fig. 26. Average light output from in-situ Far Detector strips as a function of distance from their center for normally incident MIPs. The data shown are from stopping cosmic-ray muons, for which containment criteria cause lower statistical precision at the ends of the strips.

MINERVA

- MINERvA incorporates several improvements in tracking resolution
- Triangularly extruded scintillator bars allows track position to be estimated by light-sharing fractions
- Three tracking views. Resolves ambiguity when track travels along one of the strip directions.





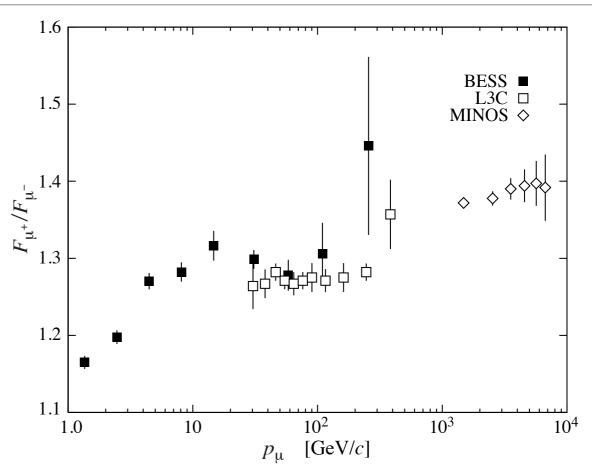


Why magnetize?

- <u>Containment:</u> A magnetic field can keep muons from exiting the sides of your detector
- Momentum measurement: If the muon does exit your detector, the curvature of the track tells you the momentum even when you couldn't otherwise get it from the range of the particle
- Charge sign: There are physics measurements in knowing the charge sign of the muons in your detector. Crucial for the "golden channel" at a neutrino factory:

$$\underbrace{\mu^+}_{} \to e^+ \bar{\nu}_{\mu} \nu_e \ \langle \nu_{\mu} \to \nu_{\mu} \nu_{\mu} \rangle \\
\nu_{e} \to \nu_{\mu}$$

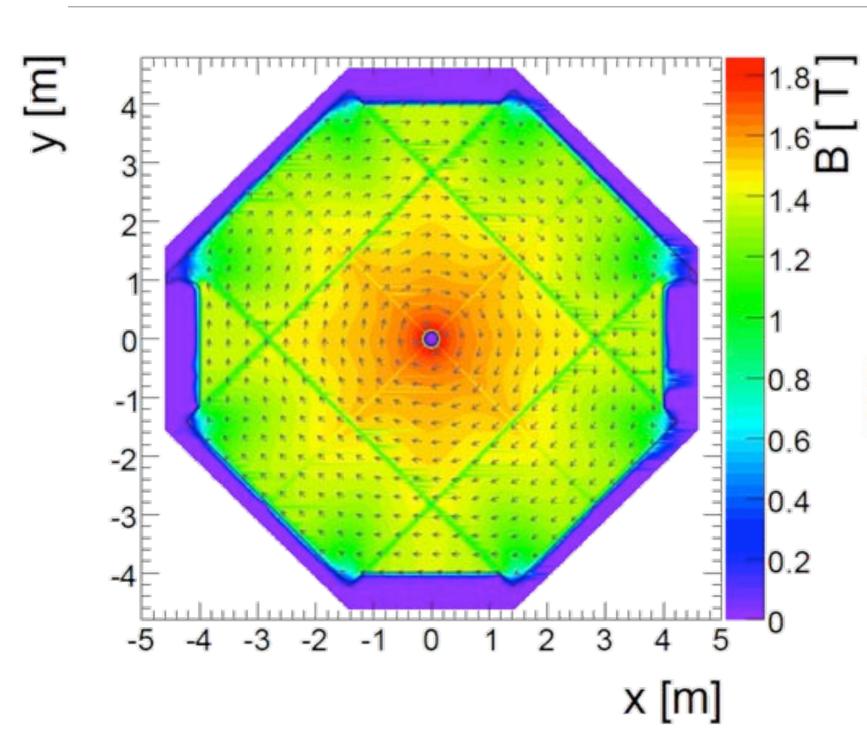
production oscillation



Cosmic-ray
$$\mu$$
+/ μ - ratio

$$ar{
u}_{\mu} + A
ightharpoonup \mu^{+} + X$$
 $u_{\mu} + A
ightharpoonup \mu^{-} + X$
 $\underline{
u}_{\mu} + X
ightharpoonup \mu^{-} + X
ightharpoonup$
"wrong" sign!

Magnetic field in MINOS



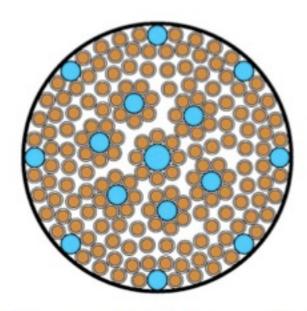
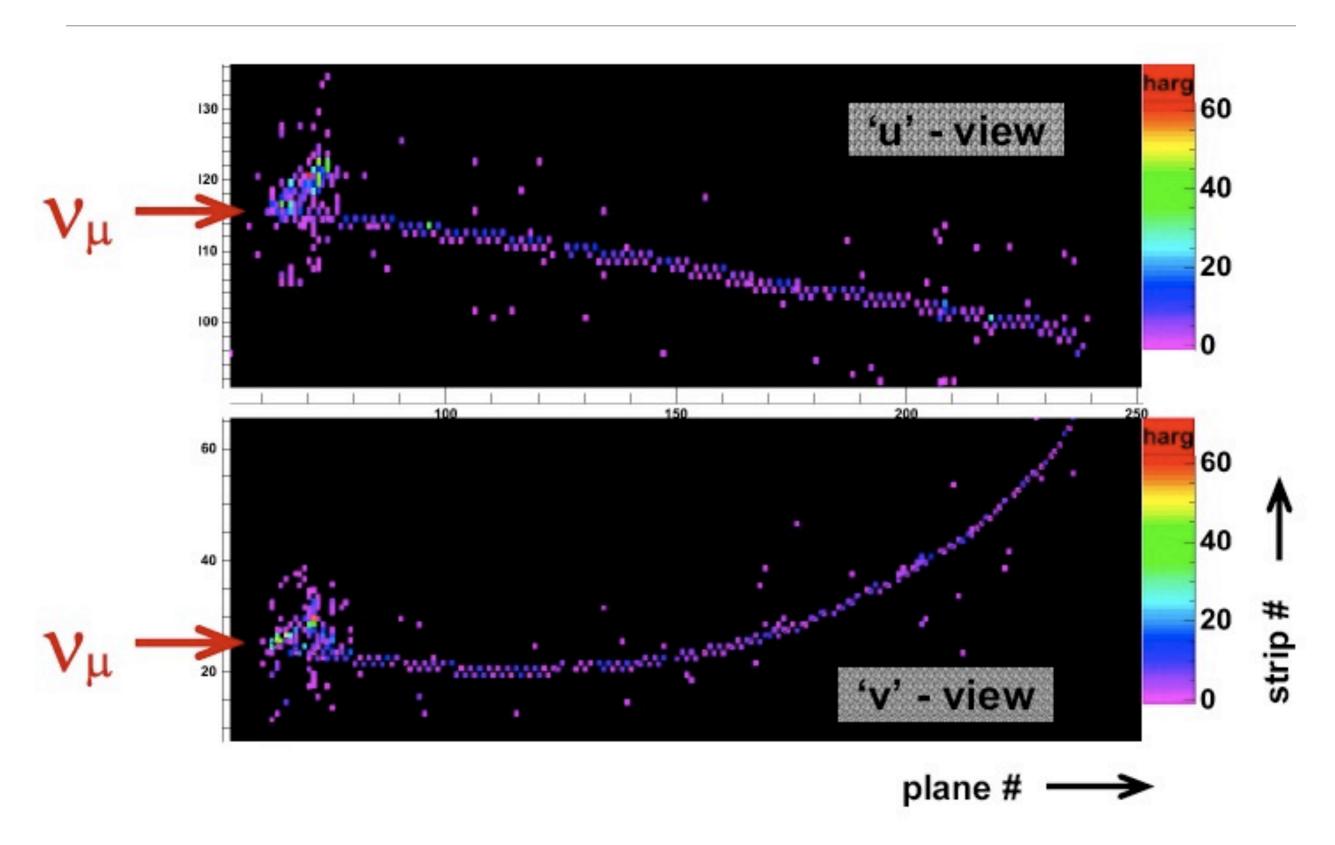


Fig. 6. Sketch of a cross section of one of the far detector supermodule coils. The larger diameter circles represent the copper cooling tubes and the smaller circles are the 190 turns of 1/0 gauge stranded copper wire. The outlines of these conductors are to-scale representations of the insulator thickness. The outer circumference of the assembly is a copper-sheet jacket that is directly cooled by eight cooling tubes.

- 15.2 kA-turn total current
- 80 A supply
- 10 gauge copper wire, water cooled

MINOS Event



Track momentum using curvature

A particle with momentum p, traveling through a constant transverse magnetic field B will travel on a circle of radius ϱ

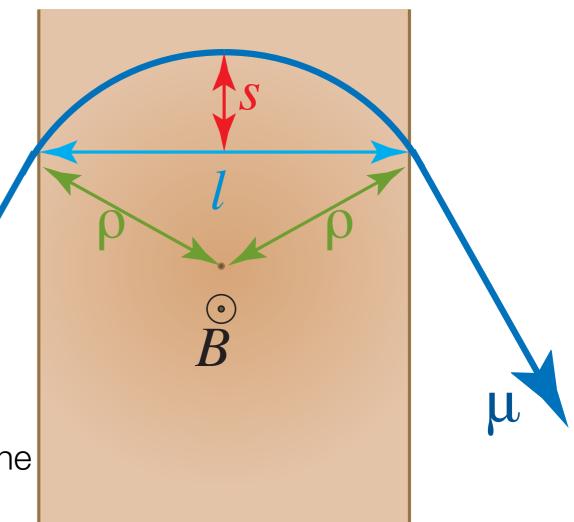
$$p[\text{GeV}/c] = 0.2998B[\text{T}]\rho[\text{m}]$$

$$\rho = \frac{l^2}{8s} + \frac{s}{2}$$

$$p \simeq 0.3 \frac{Bl^2}{8s}$$

Measurement of sagitta and chord gives you momentum. Detector resolution on sagitta is the same as the momentum resolution:

$$\left|\frac{\delta p}{p}\right| = \left|\frac{\delta s}{s}\right|$$



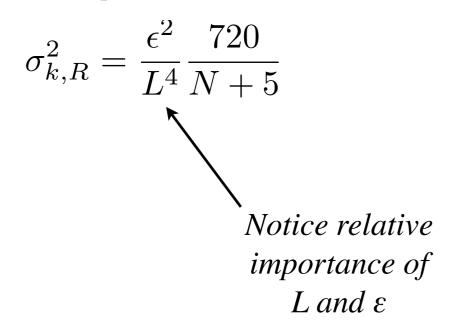
More common to talk about the track curvature

$$k = \frac{1}{\rho}$$

which has roughly Gaussian errors.

Curvature errors for multiple position samples

The uncertainty in curvature for a track which travels a distance *L* in a magnetic field *B* whose position is sampled *N* times at uniform intervals with a position uncertainty ε has been worked out by Gluckstern [NIM 24 (1963) 381-389]:



 Gluckstern has also worked out the contribution to the uncertainty in the curvature from multiplescattering:

$$\sigma_{k,M.S.}^2 = \frac{KC_N}{L}$$

• *K* is the RMS projected multiple scattering angle per unit thickness *x*

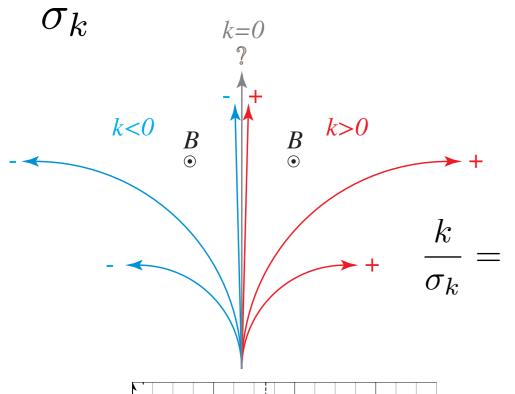
$$= \frac{\theta_0}{\sqrt{3}x}$$

$$= \frac{13.6 \text{ MeV}}{\beta cp} z \sqrt{\frac{1}{3xX_0}} \left[1 + 0.038 \ln(x/X_0) \right]$$

- C_N is a constant from lookup table. $C_N=1.43$ for large N.
- x is the distance traveled in the medium
- z is the charge of the particle

How well do we measure track curvature?

 $rac{k}{-}$ determines how well the track curvature, and hence sign is known



$$\frac{k}{\sigma_k} \sim k = \frac{0.3B}{p}$$

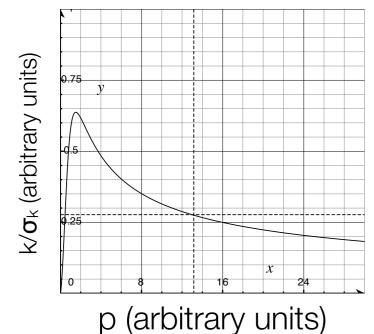
$$\sigma_k^2 = \sigma_{k,R}^2 + \sigma_{k,M.S.}^2$$

$$= \frac{1}{\left(\frac{720\epsilon^2 p^2}{L^4(N+5)} + 0.0079C_N \sqrt{\frac{p^2 + m^2}{xX_0}} (1 + 0.038 \log \frac{x}{X_0})\right)^{\frac{1}{2}}}$$

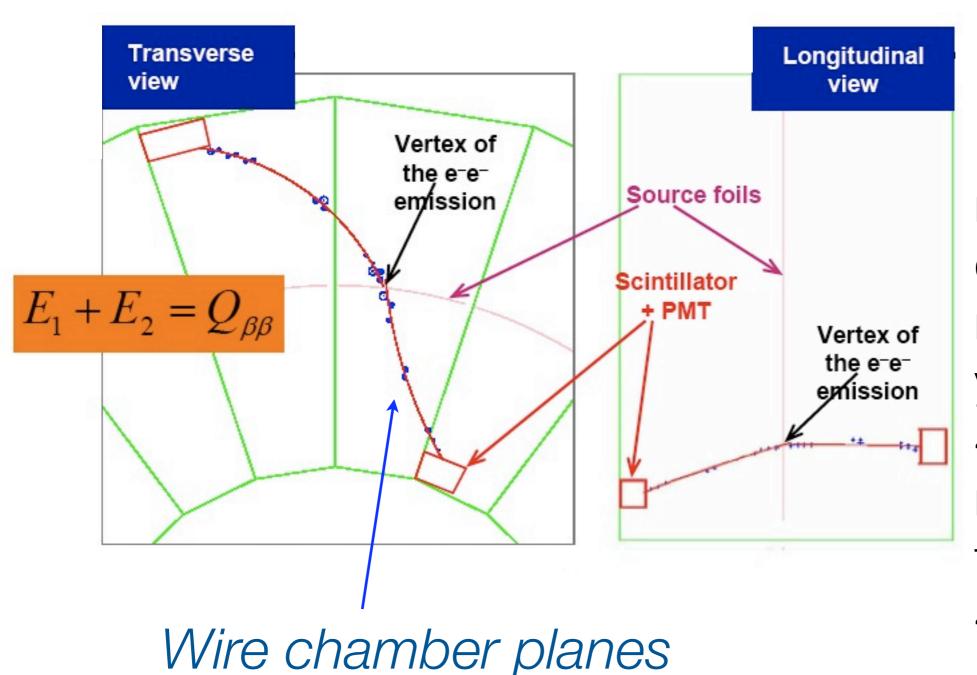
units: [T], [GeV], [m]

Remember: $L \propto p$, $N \propto p$, and $x \propto p$

- High field
- Small ε
- Large L (low Z to keep dE/dx low and range high)
- Large X₀ (low Z)
- "Just" right momentum (see plot at left)
- MINOS: 10% momentum resolution using curvature



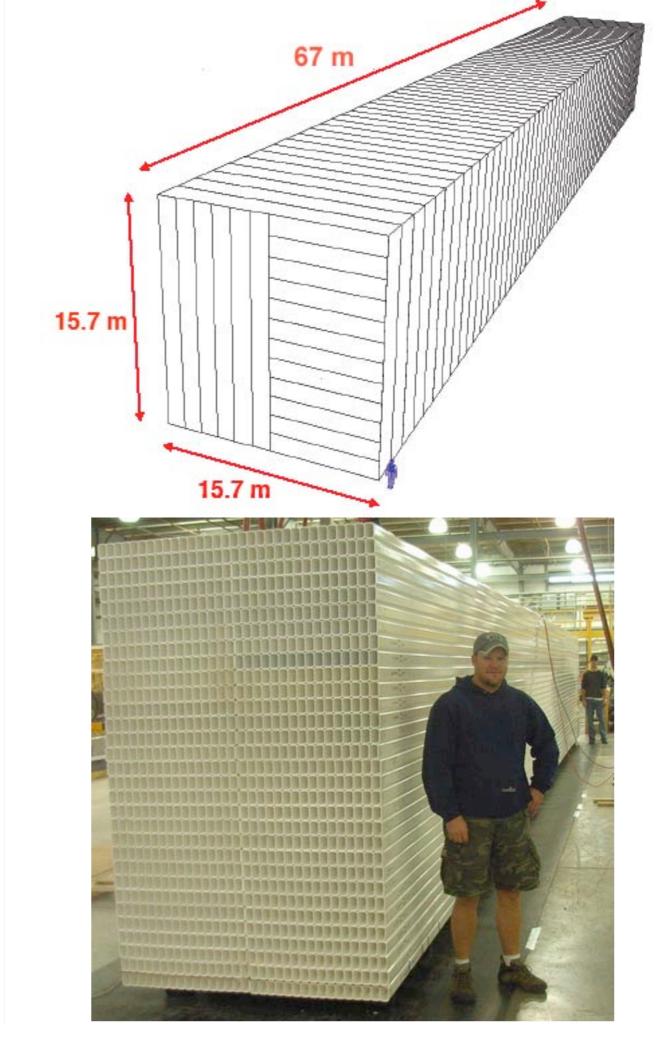
Tracking in the NEMO-3 detector $(2\nu\beta\beta)$



Low density medium and excellent sagitta measurement yields about a 4% measurement for electrons at 4 MeV

The NOvA Experiment

- NOvA is a second generation experiment on the NuMI beamline which is optimized for the detection of v_µ→v_e and v̄_µ→v̄_e oscillations
- NOvA is:
 - An upgrade of the NuMI beam intensity from 400 kW to 700 kW
 - A 15 kt "totally active" tracking liquid scintillator calorimeter sited 14 mrad off the NuMI beam axis at a distance of 810 km
 - A 215 ton near detector identical to the far detector sited 14 mrad off the NuMI beam axis at a distance of 1 km



Liquid scintillator

(14.8M liters, 12.6 ktons)

Contained in 3.9 x 6.6 cell cells of length 15.7 m

- -18 m attenuation length
- -5.5% pseudocumene

Extruded PVC

(5.4 ktons)

15% anatase TiO₂ for high reflectivity

Wavelength shifting fiber

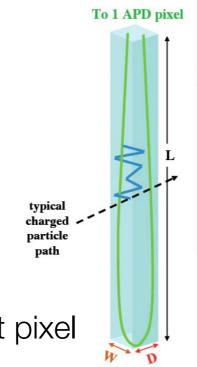
(18k km)

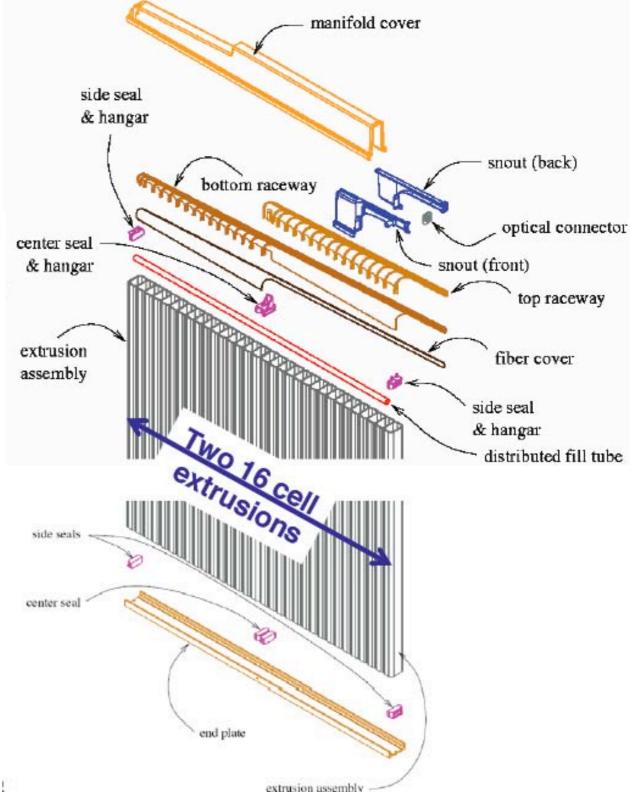
- 0.7 mm diameter
- Looped, both ends to same readout pixel

Avalanche photodiodes (APD)

(14k boards, 32 channels each)

- 85% quantum efficiency at long wavelengths
- Collect 30 photoelectrons per muon crossing at far end of cell



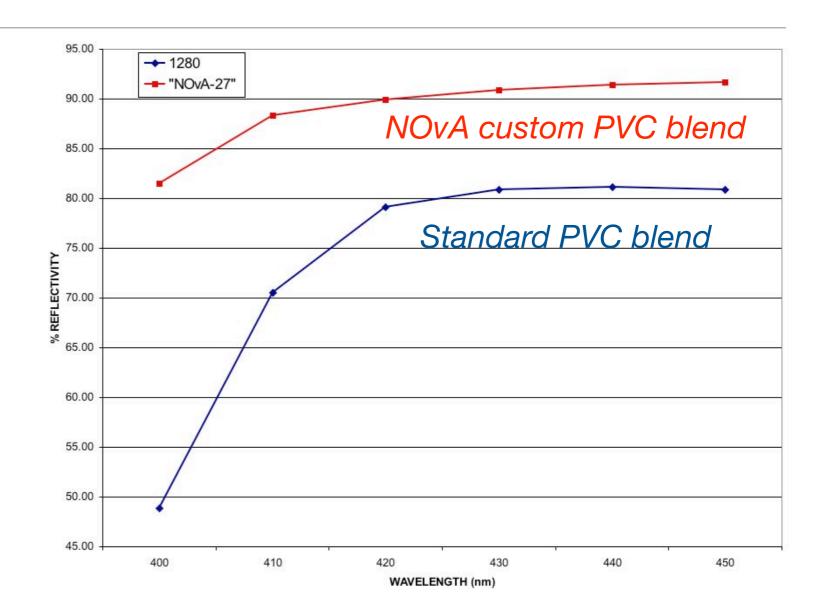


Detector design

Wall reflectivity

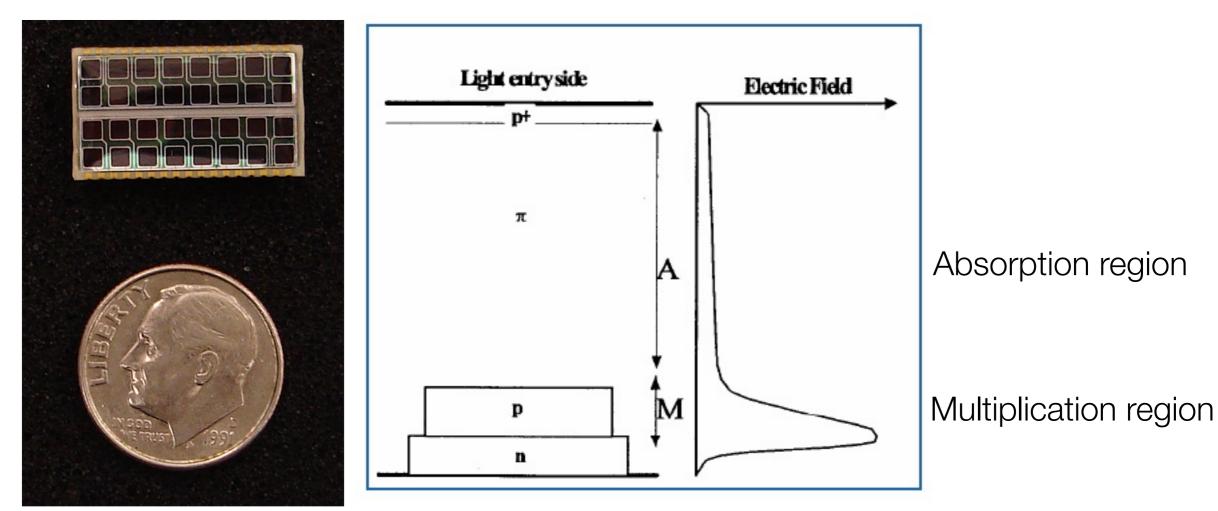
- In NOvA cell, a
 photon typically
 bounces off the cell
 walls 10 times before
 being captured by a
 fiber
- This makes the reflectivity of the cell wall of crucial importance to maximizing light output:
 - \triangleright 0.8¹⁰ = 0.11
 - $\triangleright 0.9^{10} = 0.35$

10% improvement in reflectivity yields factor 3 more light!



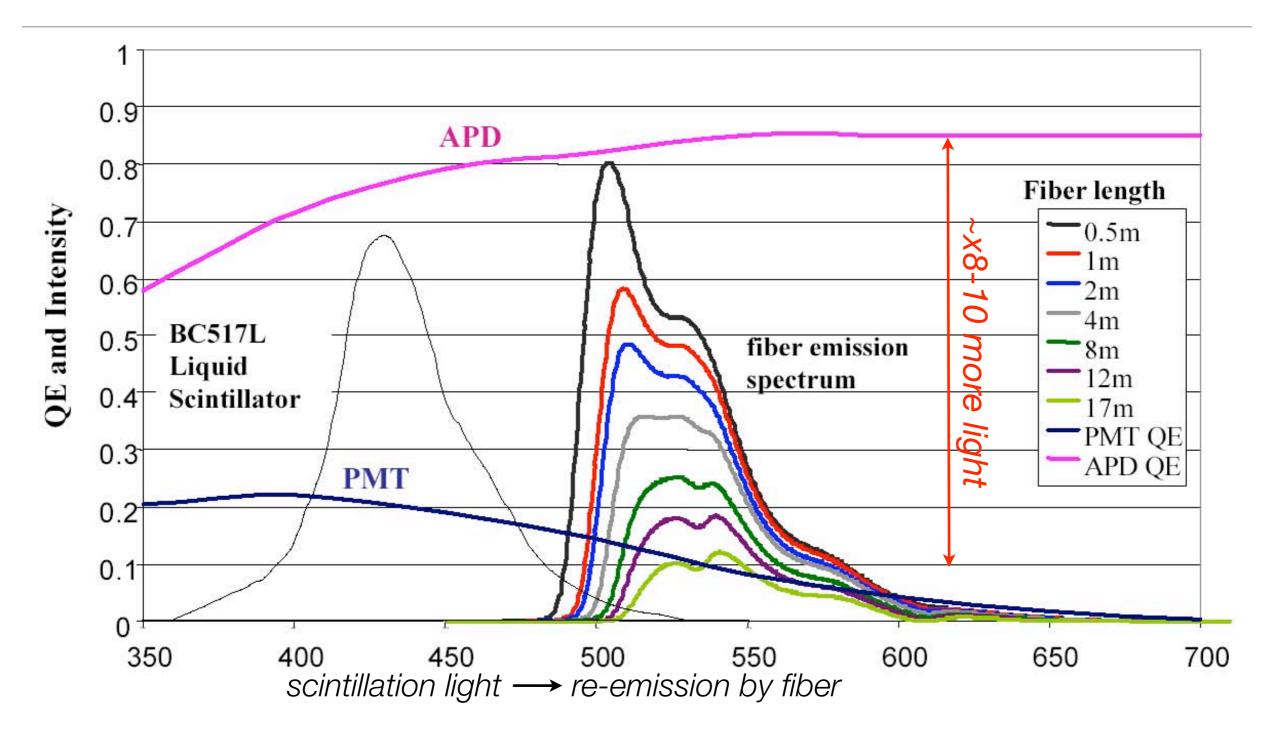
Wall reflectivity is issue for other scintillator detectors which co-extrude scintillator with a TiO2 reflective coating

Avalanche photo diodes (APD)

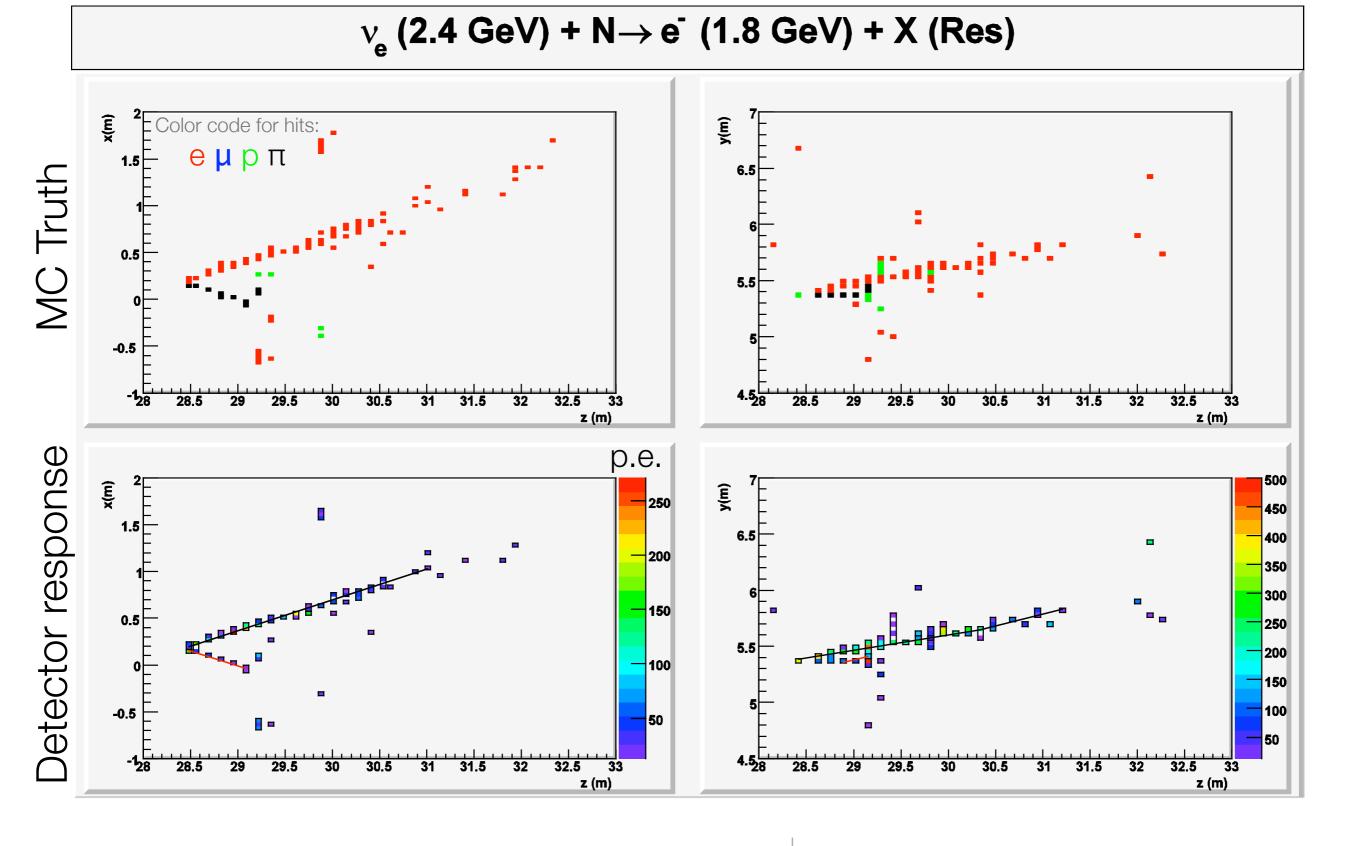


High (80%) quantum efficiency even into UV Large dark currents - must be cooled to -15°C to get noise down to ~10 pe equivalent Low gains, x100

NOvA Fiber and Photodetector



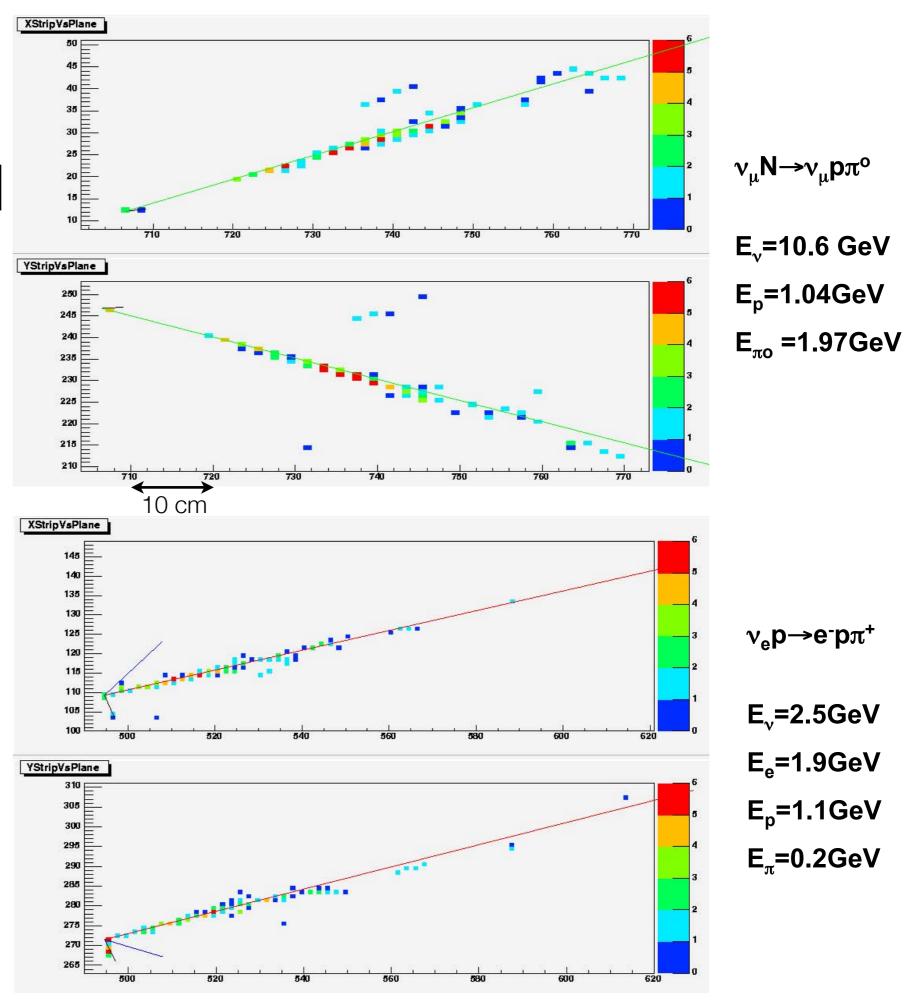
The high QE of APD's, especially at long wavelength, is crucial to NOvA performance



Electron neutrino signal event

Electron and pion tracks reconstructed

Sample signal and background events in NOvA



	Neutrino Running	Antinetrino Running	Total	Efficiency (Includes fiducial cut)
v_e signal	75.0	29.0	104	36%
Backgrounds:	14.4	7.6	22	
ν _μ NC	6.0	3.6	9.6	0.23%
ν _μ <i>CC</i>	0.05	0.48	0.53	0.004%
Beam v _e	8.4	3.4	11.8	14%
FOM	19.8	10.5	22.1	

Numbers generated assuming:

$$\sin^2(2\theta_{13}) = 0.10$$
, $\sin^2(2\theta_{23}) = 1.0$, and $\Delta m_{32}^2 = 0.0024$ eV²

Event selection

Calculations based on $\sin^2 2\theta_{13} = 0.1$ with matter effects turned off. 2 GeV NBB beam.

